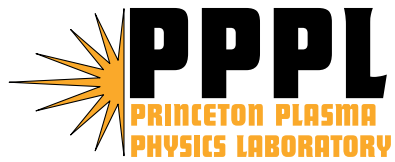


# Neutral Transport Simulations of Gas Puff Imaging Experiments on Alcator C-Mod

D. P. Stotler, B. LaBombard<sup>1</sup>,  
J. L. Terry<sup>1</sup>, and S. J. Zweben

Princeton Plasma Physics Laboratory  
Princeton University  
Princeton, NJ 08543

<sup>1</sup>MIT Plasma Science and Fusion Center  
Cambridge, MA 02139



Note: This poster is available on the Web at:  
<http://w3.pppl.gov/degas2/>

# INTRODUCTION

- Tokamak edge ideal for comprehensive study of turbulence,
  - Accessible with probes
    - ⇒ directly measure  $n_e$ ,  $T_e$ , and other properties.
  - Relatively low  $T_e$  facilitates use of atomic physics as basis for diagnostics.
  - Potential payoff great because edge sets boundary conditions for core transport,
    - \* E.g., internal transport barriers, H-mode pedestal.
- Gas Puff Imaging (GPI) experiments designed to measure 2-D structure of edge turbulence,
  - Compare with 3-D nonlinear simulations.
  - And with turbulence measured by probes,
  - Puff neutral gas (e.g.,  $D_2$ ) near outer wall,
    - \* View with fast, high res. camera light from electron impact excitation of gas,
    - \* Use sightline  $\parallel \vec{B}$  to see radial & poloidal structure,
- Explore relation between images & plasma fluctuations with DEGAS 2 neutral transport code,
  - Straightforward because puff does not perturb plasma,
  - Emitted light brighter than background,
  - Material surface interactions should not be important.
- Experimental presentation: O-02 J. L. Terry et al.

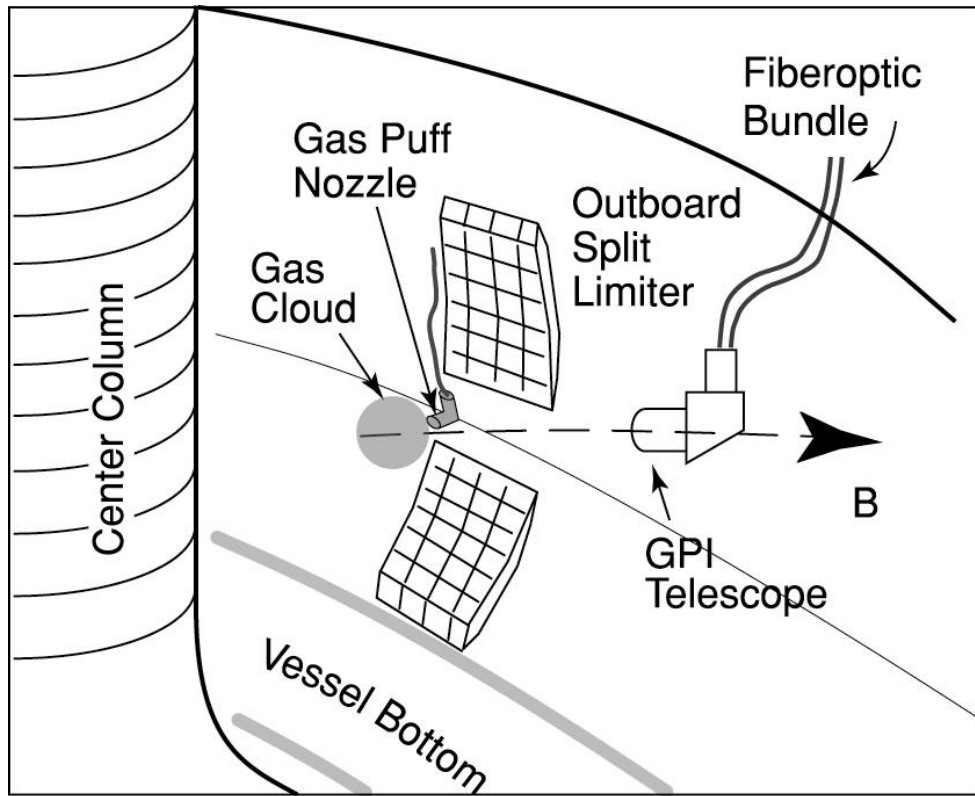


Fig. 1 - Zweben APS '01

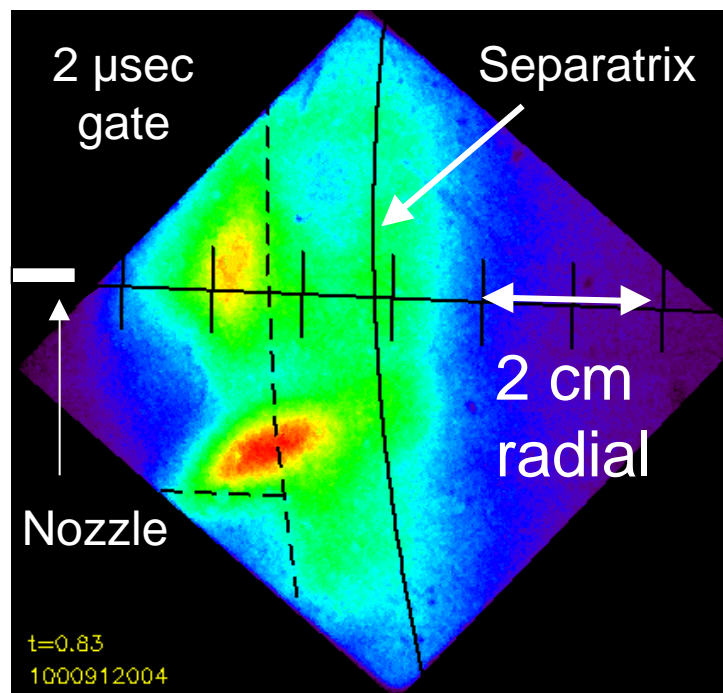
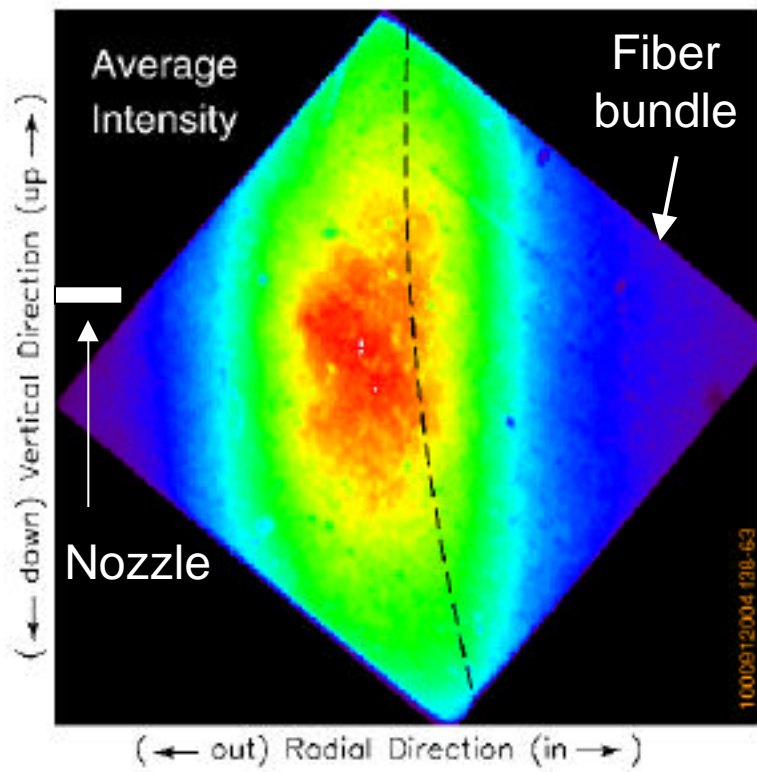


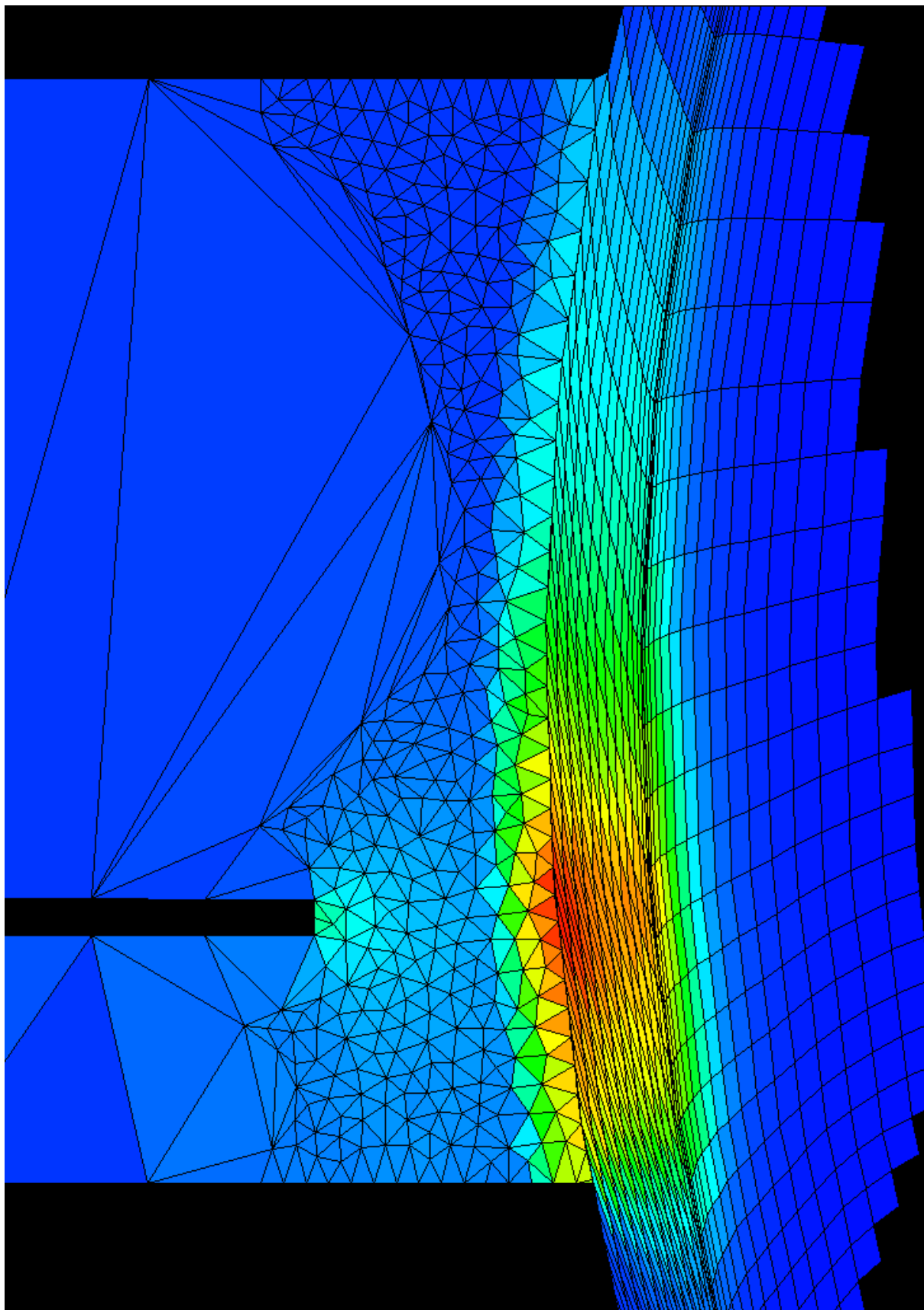
Fig. 2 - Zweben / APS '01

# DESCRIPTION OF DEGAS 2 SIMULATIONS

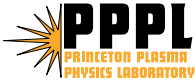
- Alcator C-Mod Geometry:
  - Start with outline of vacuum vessel,
    - \* Including gas puff nozzle & surrounding structures.
  - EFIT equilibrium for time of interest  $\Rightarrow$ 
    - \* 2-D plasma mesh set up using DG & Carre,
    - \* Bunch surfaces & grid points to get resolution 3 mm or smaller in region of interest.
  - Divide puff region into  $\sim 3$  mm triangles using Triangle.
- Simulations 2-D axisymmetric for now,
  - Output is averaged over toroidal angle.
  - $\Rightarrow$  poloidal plane variation of photon emission rates.
  - Plan to add toroidal resolution  $\Rightarrow$ 
    - \* Can directly simulate fast camera views,
    - \* Quantitative comparison of image intensity,
    - \* Evaluate toroidal spatial averaging.



# DEGAS 2 Geometry for C-Mod Shot 1010622

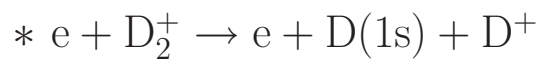
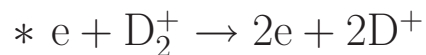
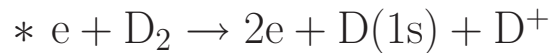
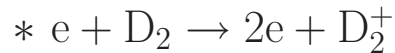
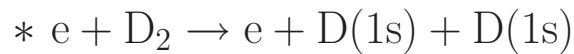


- Simulations assume steady-state.
  - Compare time scales:
    - \* Autocorrelation time for turbulence  
 $= 10 - 20 \mu\text{s}$ ,
    - \* Time for 3 eV D to travel across cloud  
 $= 1 \mu\text{s}$  (2 cm),
    - \* Timescale for emission of  $D_\alpha$  photon  
 $= 1/A_{3 \rightarrow 2} = 0.02 \mu\text{s}$ ,
    - \* Note that camera exposure times  
 $= 2 \mu\text{s}$  (60 frame/s) or  
 $4 \mu\text{s}$  ( $5 \times 10^6$  frames / s),
    - \*  $\Rightarrow$  assumption of stationary plasma OK.



- Physics:

- $D_2$ ,  $D_2^+$  dissociation, including



- $D + D^+$  elastic scattering (i.e., charge exchange),

- $D_2 + D^+$  elastic scattering,

- $e + D$  ionization,

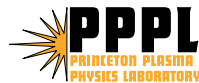
- \* “Multi-step”, i.e., collisional-radiative model.

- Neutral-neutral collisions *not* included,

- \* May not be negligible,

- \* Need realistic neutral density to treat,

- \* Can only be computed in 3-D.





– Emission rate ( $\text{m}^{-3} \text{s}^{-1}$ ) written as:

$$S_{D_\alpha} = \sum_{j=D, D_2, D_2^+} n_j f_j(n_e, T_e),$$

\* Where  $n_j$  = ground state atom & molecule density,

$$f_D = \frac{n_D(n=3)}{n_D(n=1)} A_{3 \rightarrow 2},$$

\*  $[n_D(n=3)/n_D(n=1)](n_e, T_e)$  from CR model,

\* Largely determines  $n_e, T_e$  dependence of  $f_D$ .

$$f_{D_2}, f_{D_2^+} = n_e \sum_k \langle \sigma v \rangle_k(T_e),$$

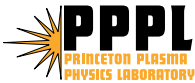
\*  $k$  = reactions leading to  $n=3$ .

– All puffs are 300 K with cosine distribution,

\* Examined sensitivity in preliminary runs,

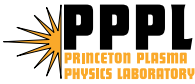
\* Run with  $(\cos \theta)^4$  distribution,

\* One with 150 K puff.

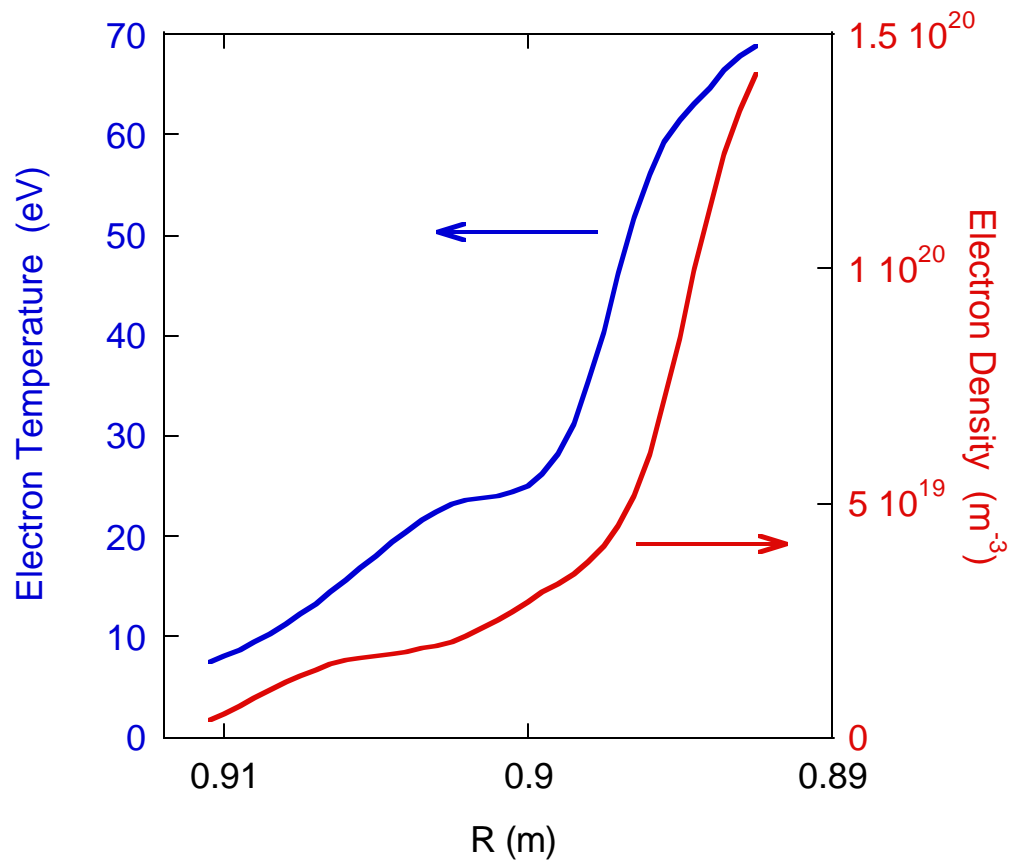


- Plasma profiles:

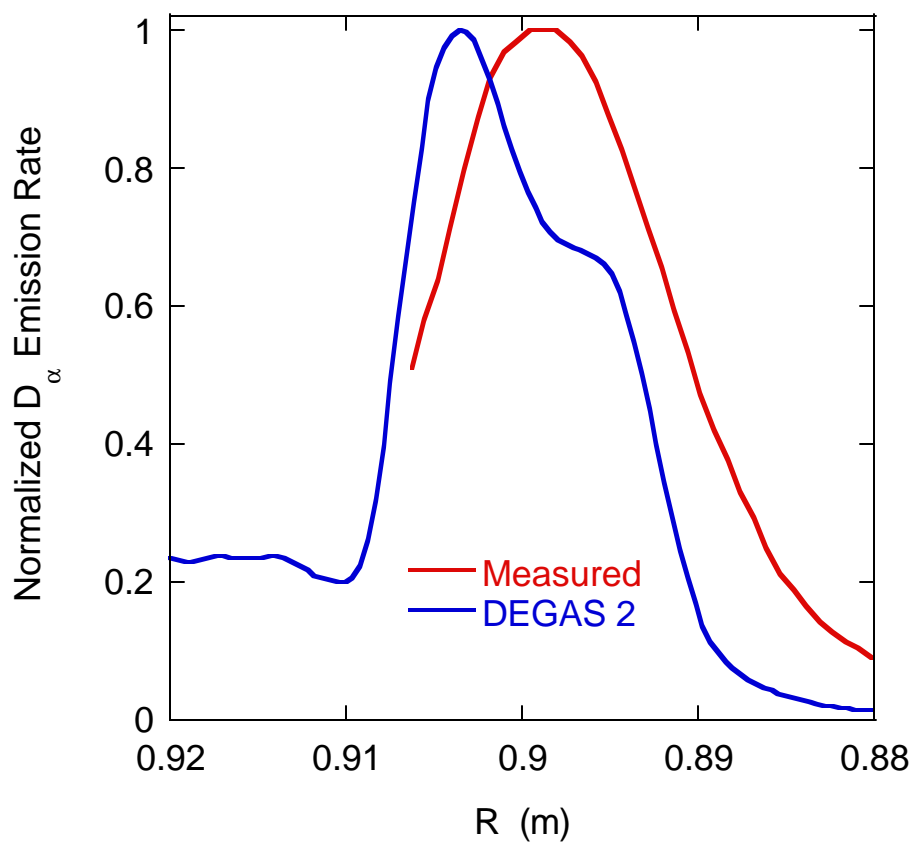
- All are taken from measured data mapped to midplane,
- Assume constant on a flux surface,
  - \* In triangulated region, estimate  $\rho$  = distance between zone center & nearest flux surface mesh zone.
- Assume  $n_i = n_e$ ,  $T_i = T_e$ .



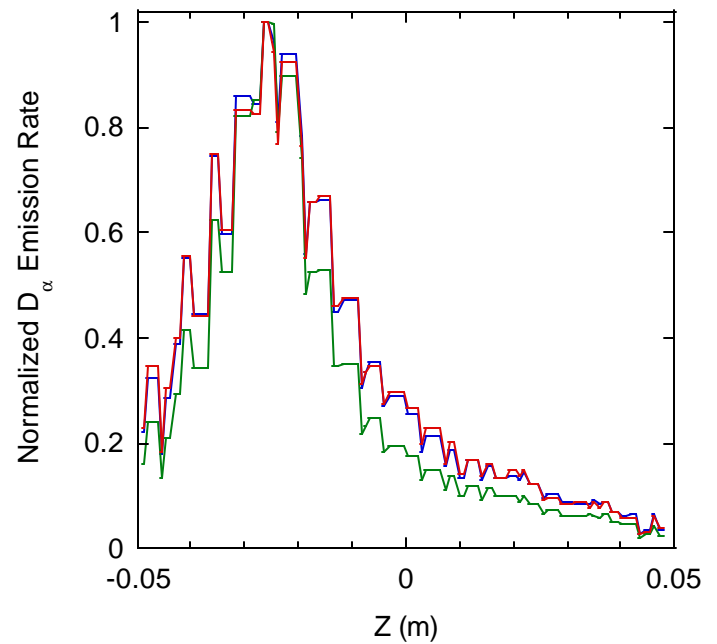
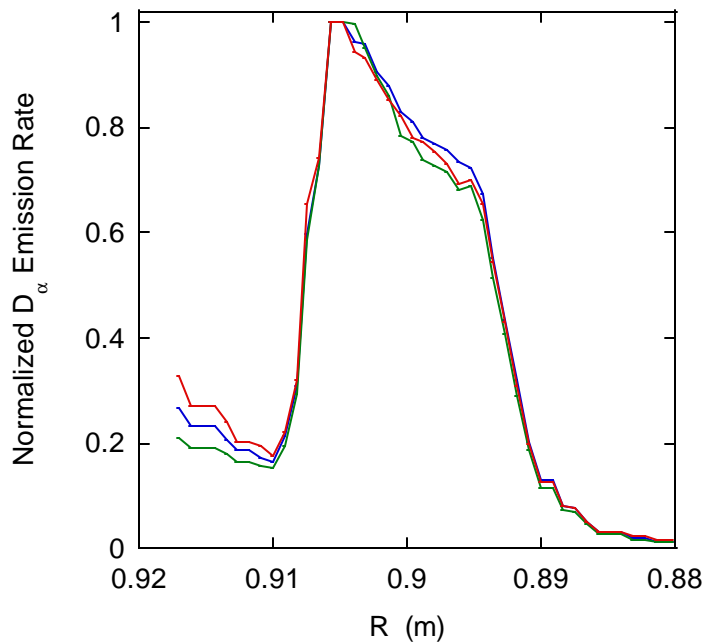
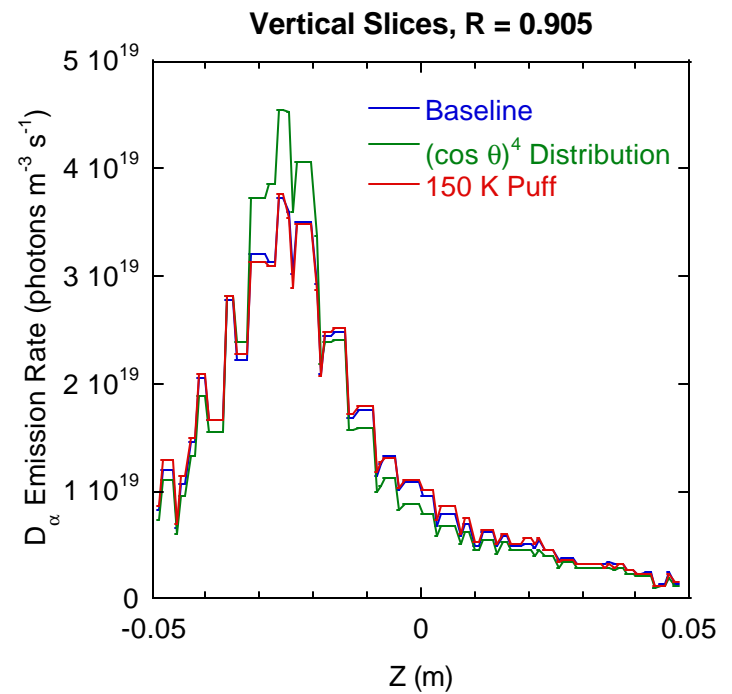
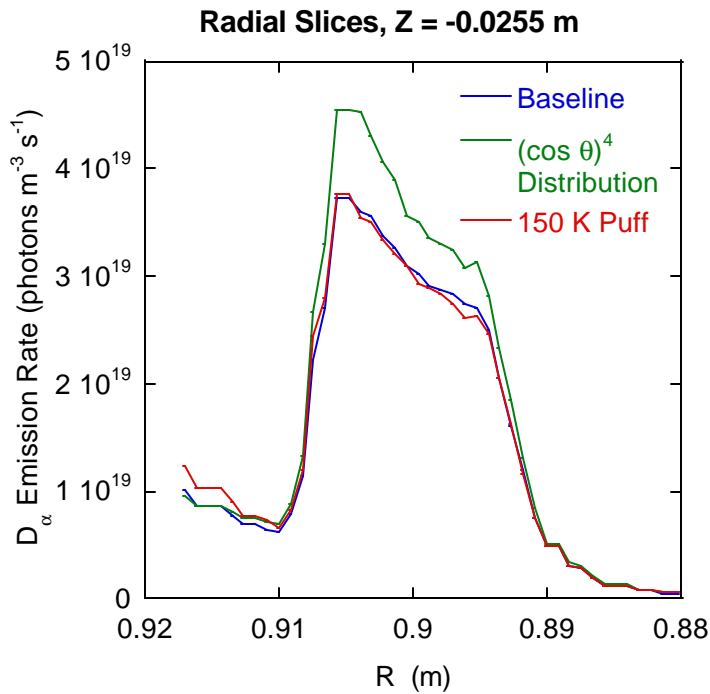
**Scanning Probe Data from C-Mod**  
**Shot 1010622006, 700 ms**



**Compare DEGAS 2 Result with Experimental Data**  
**Radial Slice at  $Z = -0.034$  m**



# Peak Location & Width of Simulated Emission Insensitive to Details of $D_2$ Distribution

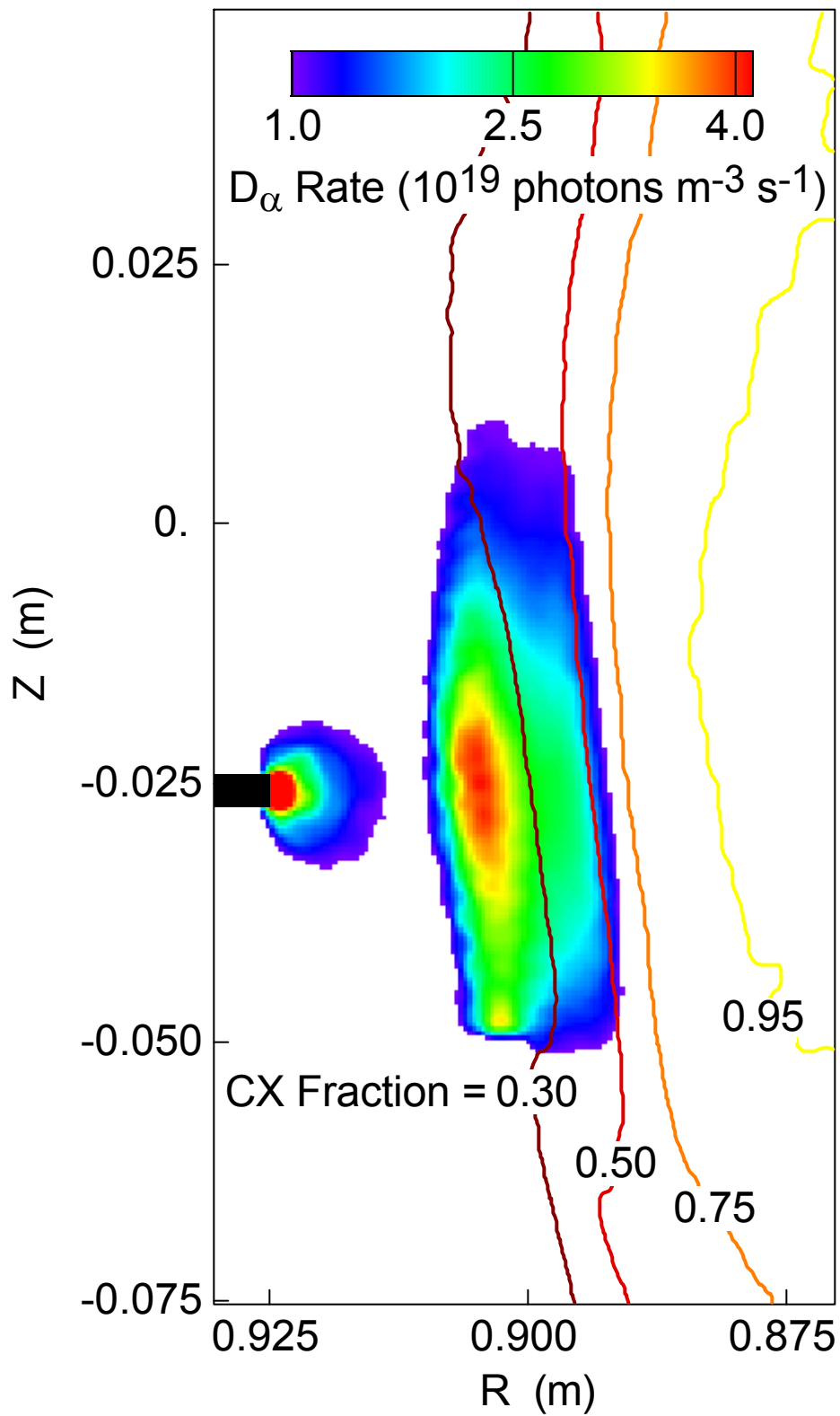


⇒ Vertical extent can be affected

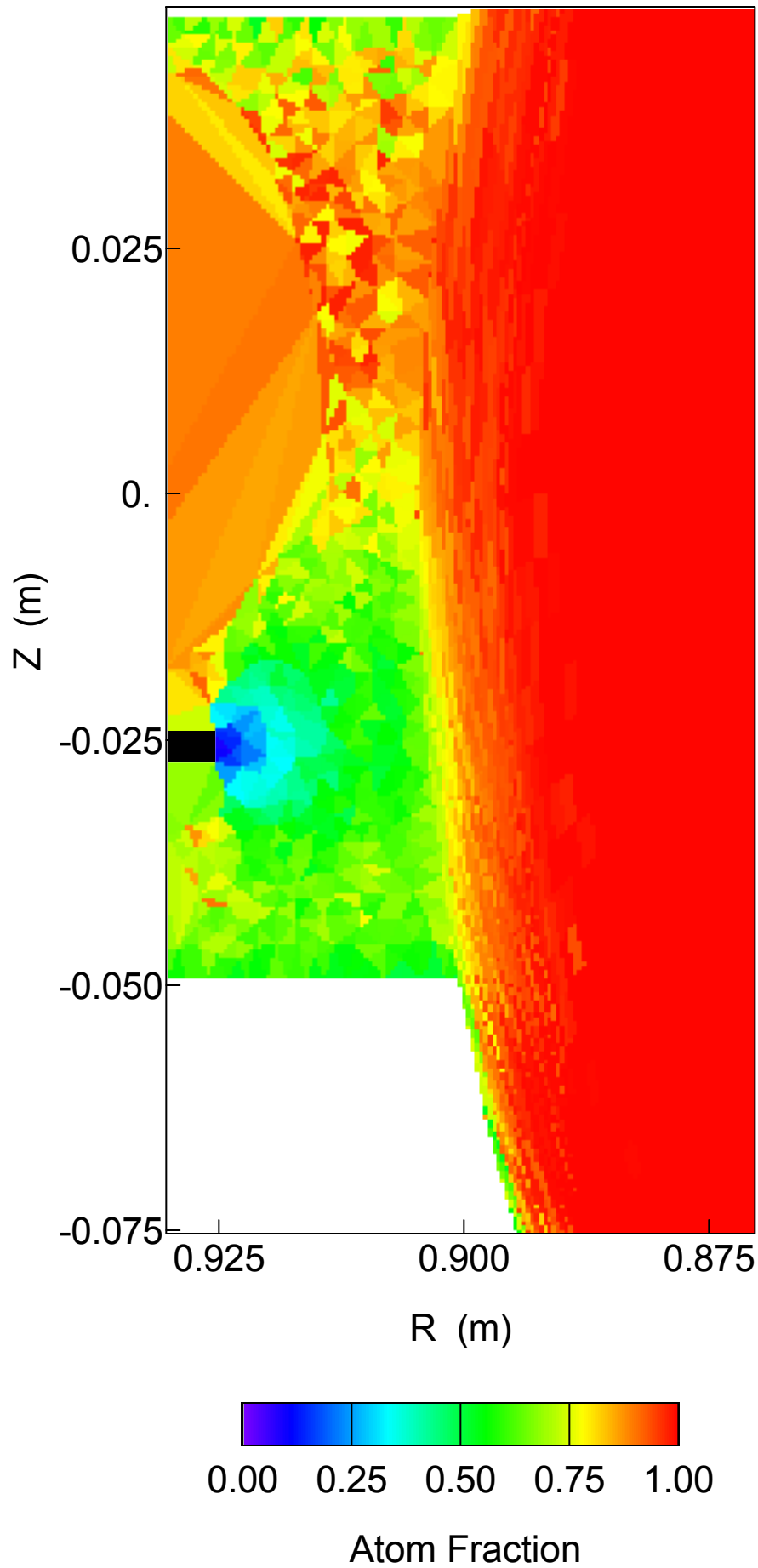
## C-MOD RESULTS

- Alcator C-Mod shot 1010622006 at 700 ms.
- Baseline computed with time-average plasma profiles,
  - 10 – 20% of atoms in cloud undergone reflection,
  - “CX fraction” have had a CX,
  - Rest from dissociation  $\Rightarrow$  ballistic trajectories.
  - $\Rightarrow \sim 50 - 65\%$  of D emission
- At peak, molecular  $D_\alpha$ s contribute  $\sim 40\%$ ,
  - $< 10\%$  for  $R \lesssim 0.9$  m.
- Compare with time-average experimental GPI images,
  - Emission peak near nozzle not seen experimentally,
  - Probe data assumed constant for  $R > 0.91$  m,
  - Nozzle peak  $\downarrow 10^{-2}$  if  $T_e < 2.5$  eV
  - Or if  $n_e < 3.6 \times 10^{16} \text{ m}^{-3}$ ,
  - Both consistent with exponential extrapolation of probe data.

# DEGAS 2 Baseline



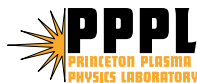
# Fraction of $D_\alpha$ Due to Atoms



- Impose 2-D perturbation on  $n_e$  and  $T_e$ ,
  - Important to understand relation between spatial variation in emission & underlying plasma fluctuations,
  - Consider ad hoc perturbation:

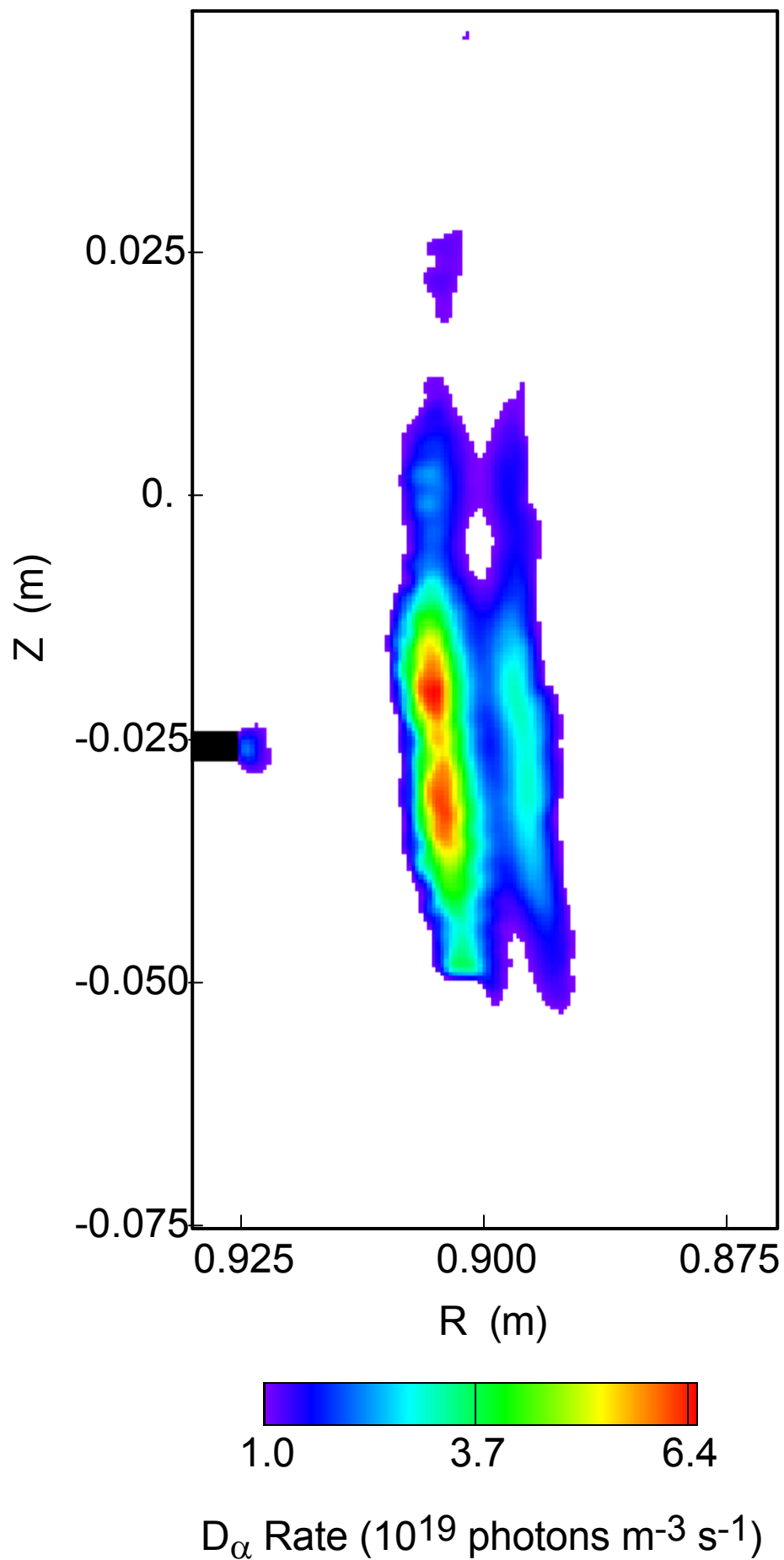
$$n'_e(R, Z) = n_e(R, Z) \left[ 1 + \frac{1}{2} \sin\left(\frac{\pi Z}{0.01}\right) \right] \times \left\{ 1 + \frac{1}{2} \sin\left[\frac{\pi(R - R_{\text{sep}} + 0.0035)}{0.005}\right] \right\},$$

- where:
  - \* The 1/2 factors make this a 50% perturbation,
    - Factor ranges from 0.25 to 2.25.
  - \* 2 cm wavelength for poloidal ( $\sim Z$ ) variation,
    - Typical size of observed emission structures.
  - \* Used only 1 cm in  $R$  because of limited radial width,
    - 0.0035 shift so innermost data point unchanged.
- Try same perturbation on  $T_e$ ,
  - \* Only difference:  $T_e$  bound between 5 and 100 eV.

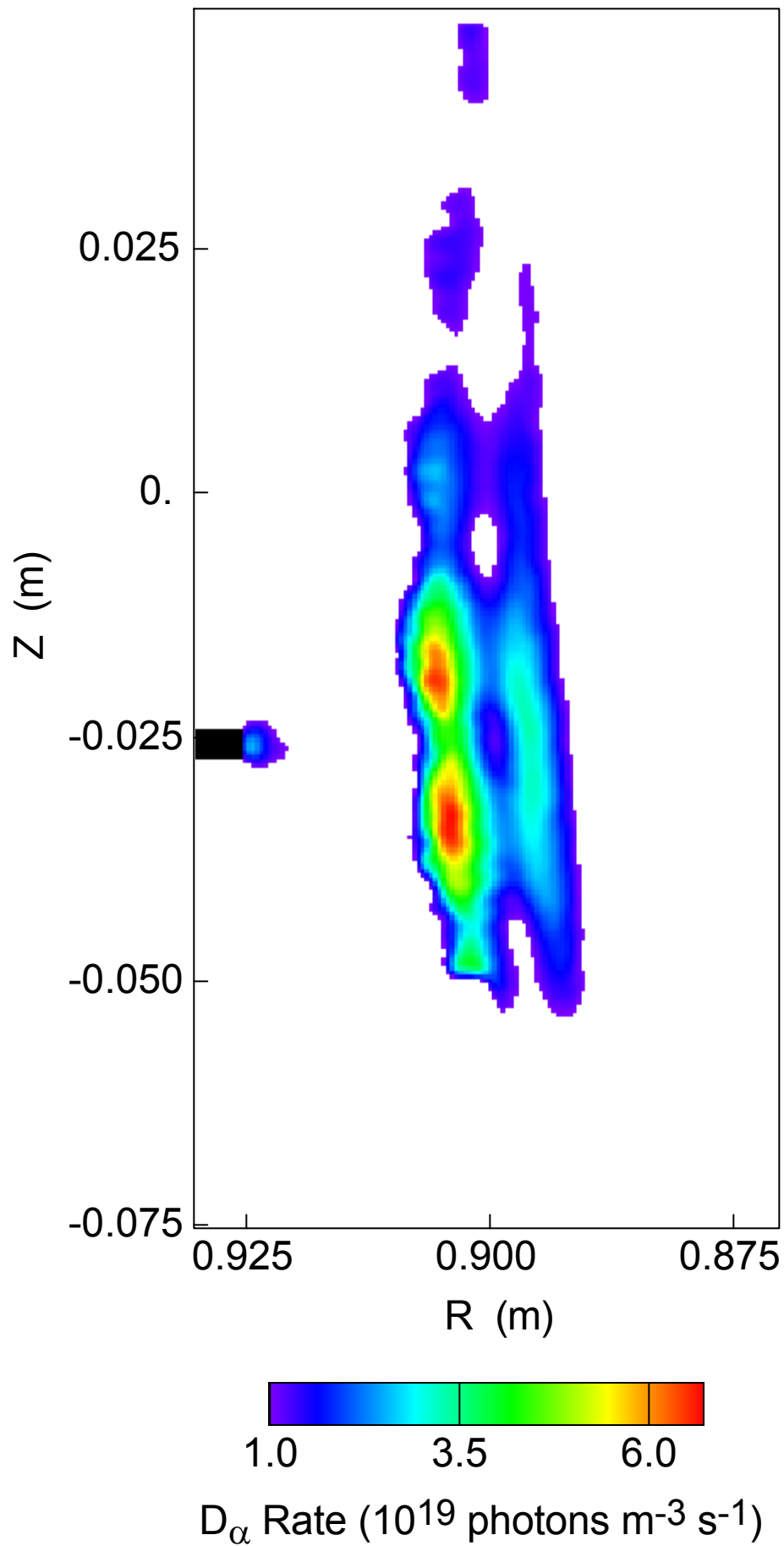




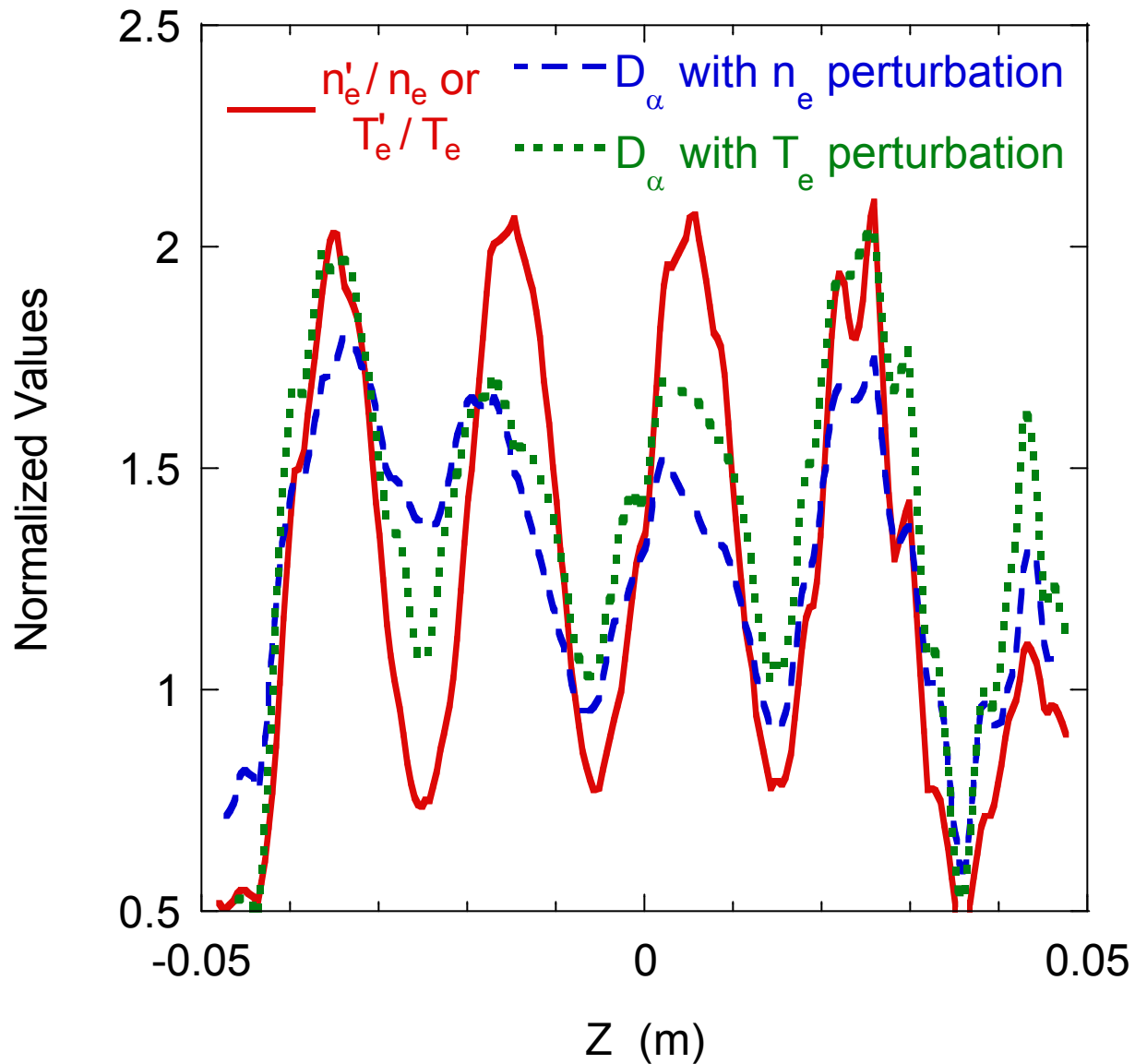
# 2-D Perturbation to Electron Density



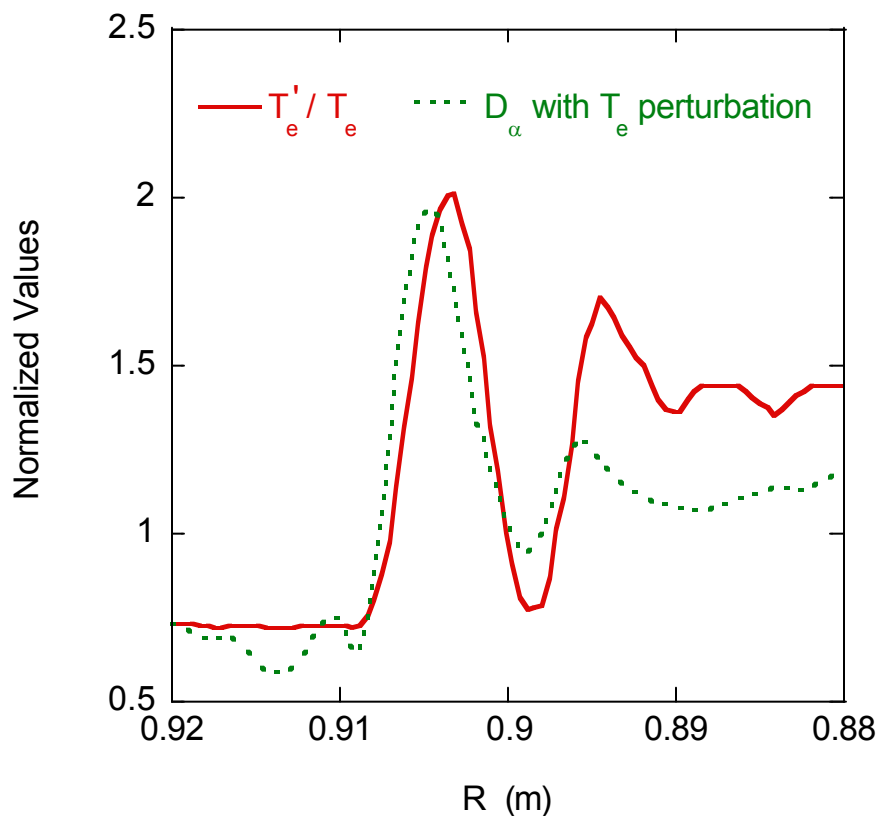
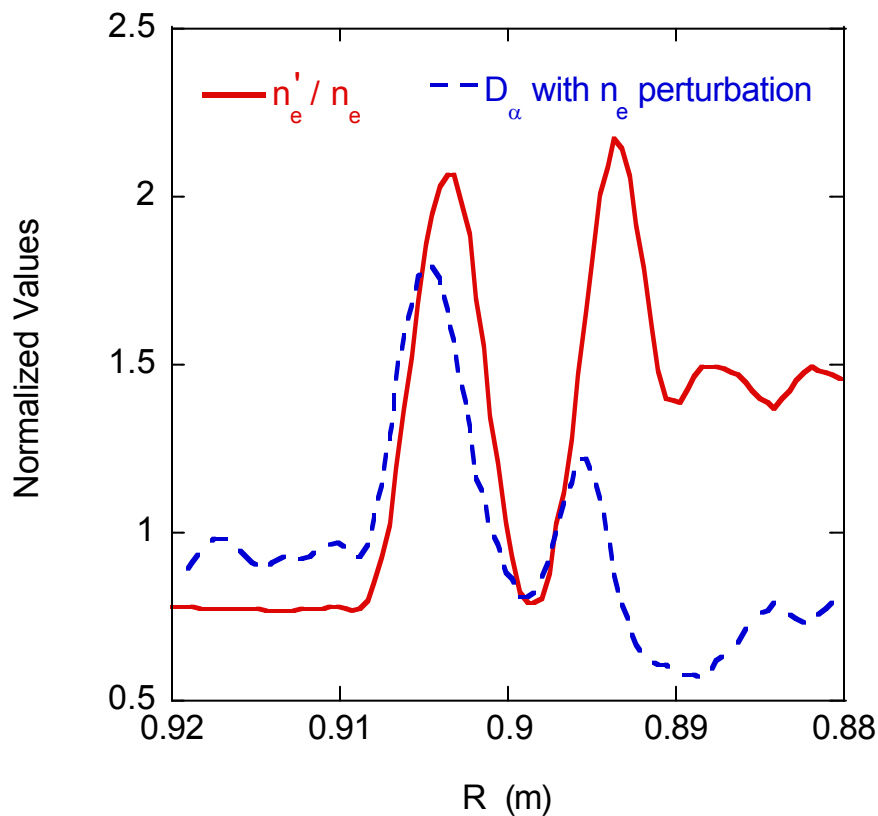
# 2-D Perturbation to Electron Temperature



# Effect of 2-D Perturbation Normalized to Unperturbed Value Vertical Slice

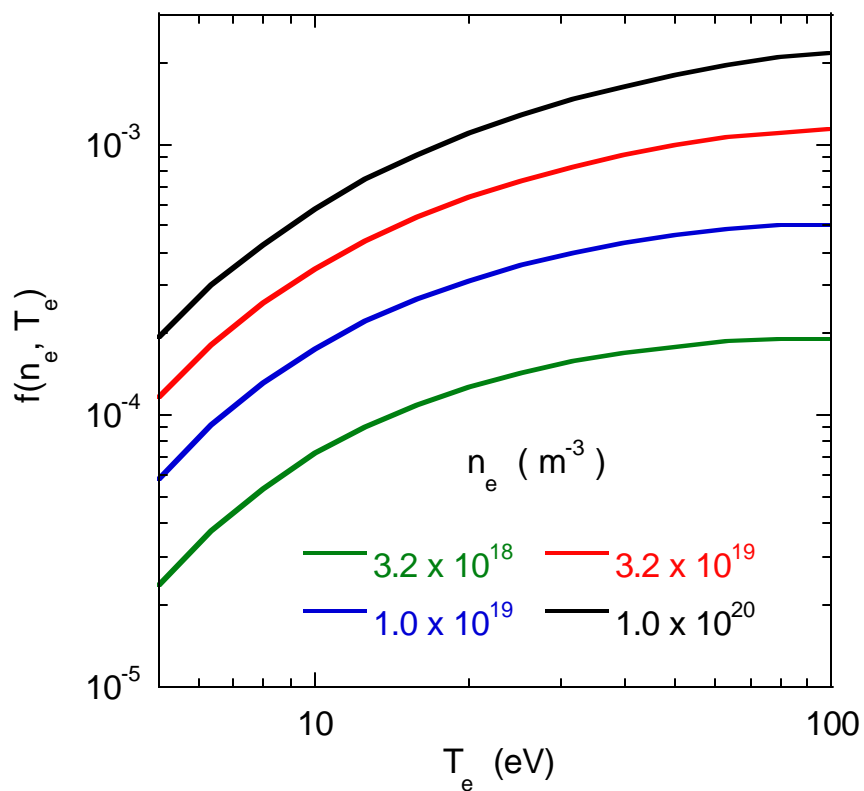


# Effect of 2-D Perturbation Normalized to Unperturbed Value Radial Slice

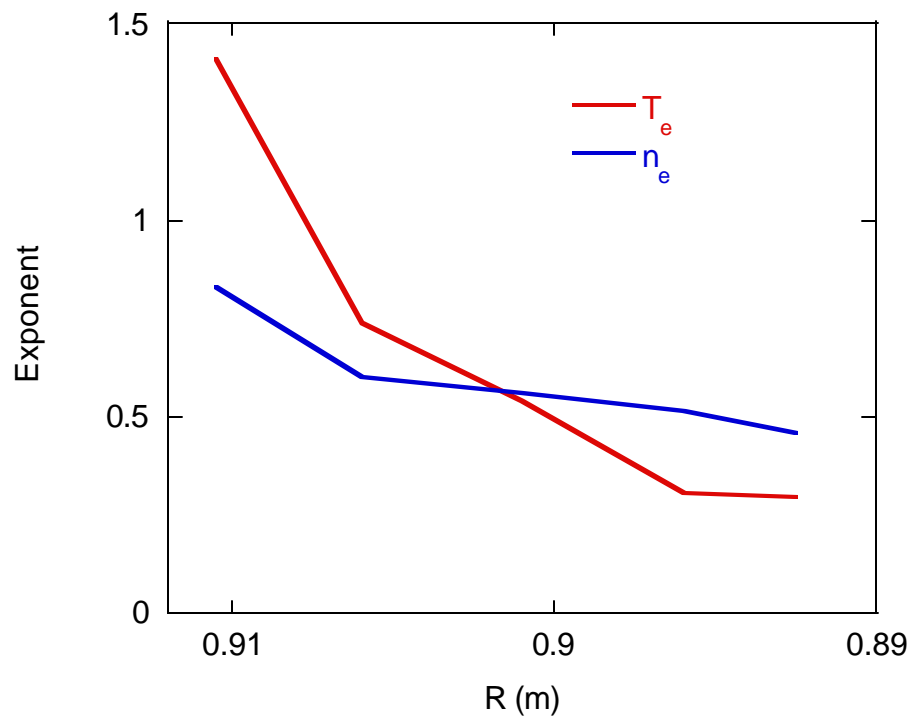


- Both simulations shows same 2-D structure,
- $\Rightarrow$  wavenumber spectrum at least similar to that of plasma turbulence,
  - \* Expect autocorrelation function & frequency spectra similar also,
  - \* Will subsequently investigate quantitatively.
- Ratio of perturbed / unperturbed emission  $\neq n'_e/n_e$  because  $\partial \ln f_D / \partial \ln n_e, \partial \ln f_D / \partial \ln T_e < 1$ .
- Further complicated by molecular contributions,
  - \*  $f_{D_2}$  and  $f_{D_2^+} \propto n_e$ ,
  - \*  $T_e$  dependence not simple,
  - \* Effective scaling varies radially.
- Simple interpretation of GPI: image patterns  $\propto n'_e/n_e$ ,
  - And insensitive to  $T_e$ ,
  - Valid only if  $n_e \lesssim 10^{18} \text{ m}^{-3}$  and  $T_e \gg 10 \text{ eV}$ ,
  - Not the case here!
  - $\Rightarrow n_e, T_e$  dependence of  $S_{D_\alpha}$  not different enough to infer perturbation amplitudes,
  - Would be simpler if  $n_e, T_e$  in phase.

**$n_e, T_e$  Dependence of  $D_a$  Emission Rate  
Contained in Ratio of  $n=3$  Density to  $n=1$**

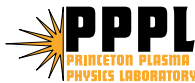


**Scaling of  $f(n_e, T_e)$  Varies  
Across Radial Profiles of 1010622**



## Shadow Fraction

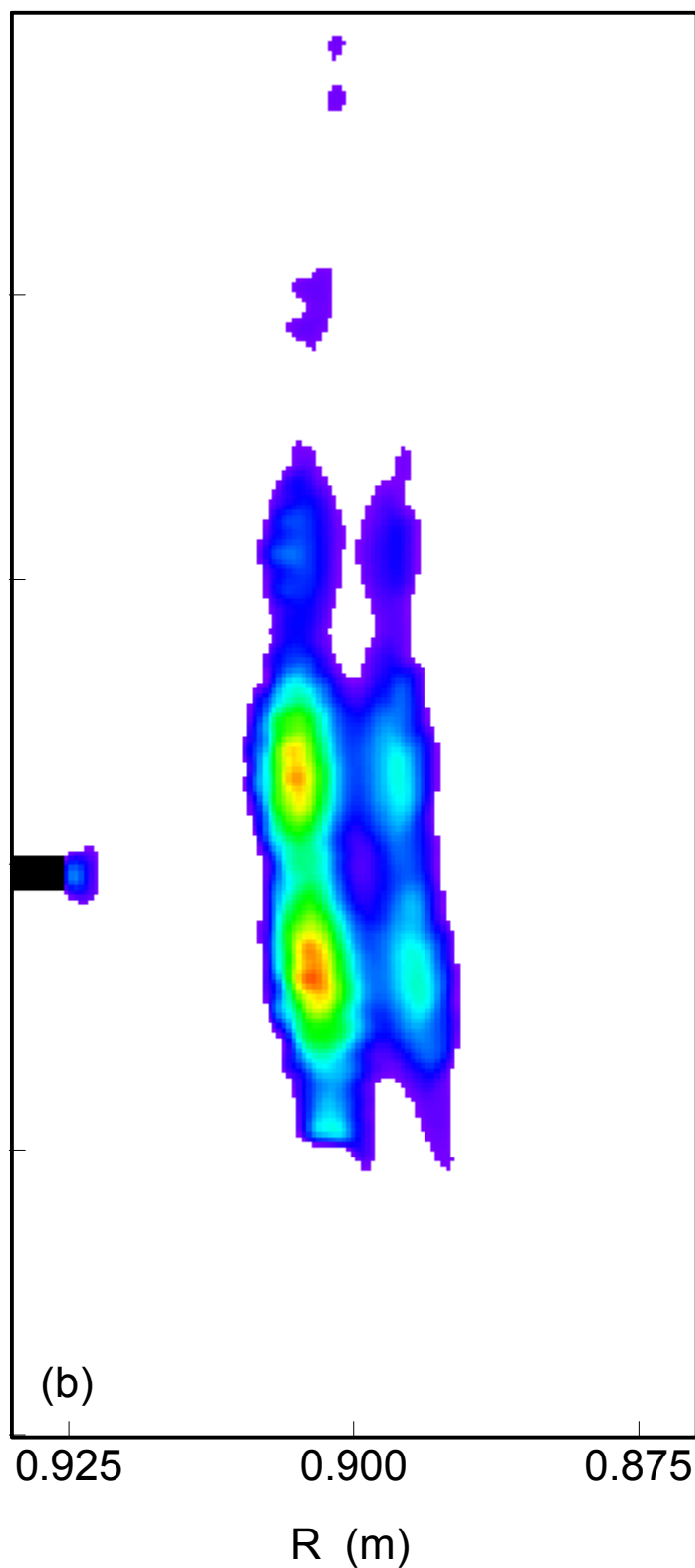
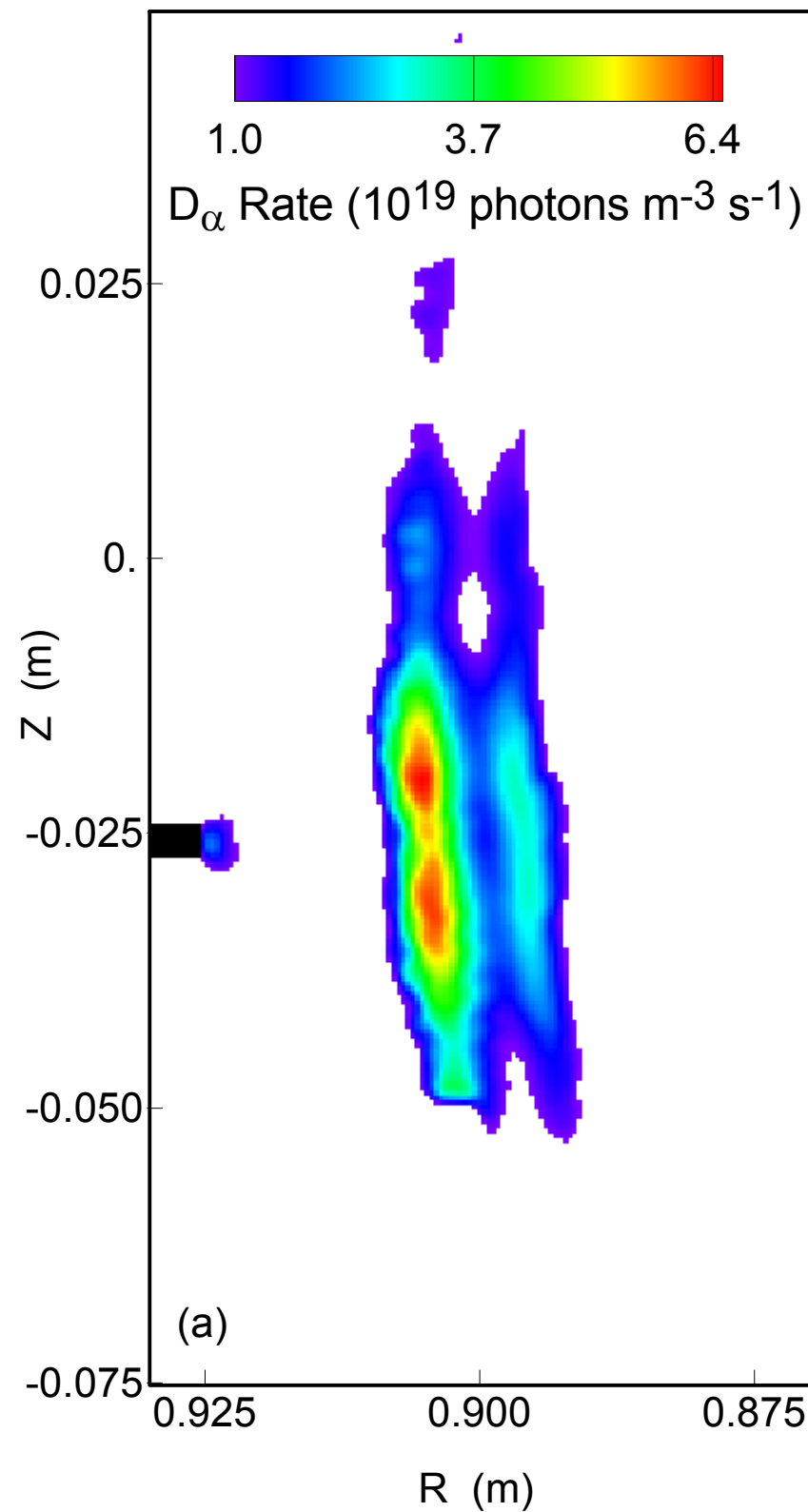
- Above focussed only on effect of perturbation on  $f_j$ ,
- They also impact  $n_j$ !
- “Shadowing effect”: ionization caused by local  $n_e$ ,  $T_e$  peak reduces light at smaller  $R$ .
- Compare images with and without shadowing,
  - “With” shadowing is as above,
  - To eliminate, use perturbed  $f_j$  and *unperturbed*  $n_j$ ,
  - “Unshadowed” clearly shows  $n_e$  perturbation structure,
  - Shadowed image smeared out,
    - \* Due to  $n_j$  reductions by  $n_e$  peaks,
    - \* And  $n_j$  increases by  $n_e$  minima.



# Runs with Electron Density Perturbation Shadowing:

with

without





- Estimate by computing:

$$F_s = \left[ \sum_j (n'_j - n_j) f'_j \right] / \sum_j n_j f_j,$$

- Where prime indicates perturbed value.
- Evaluate separately for both “perturbed” simulations.

- Structure is complicated!

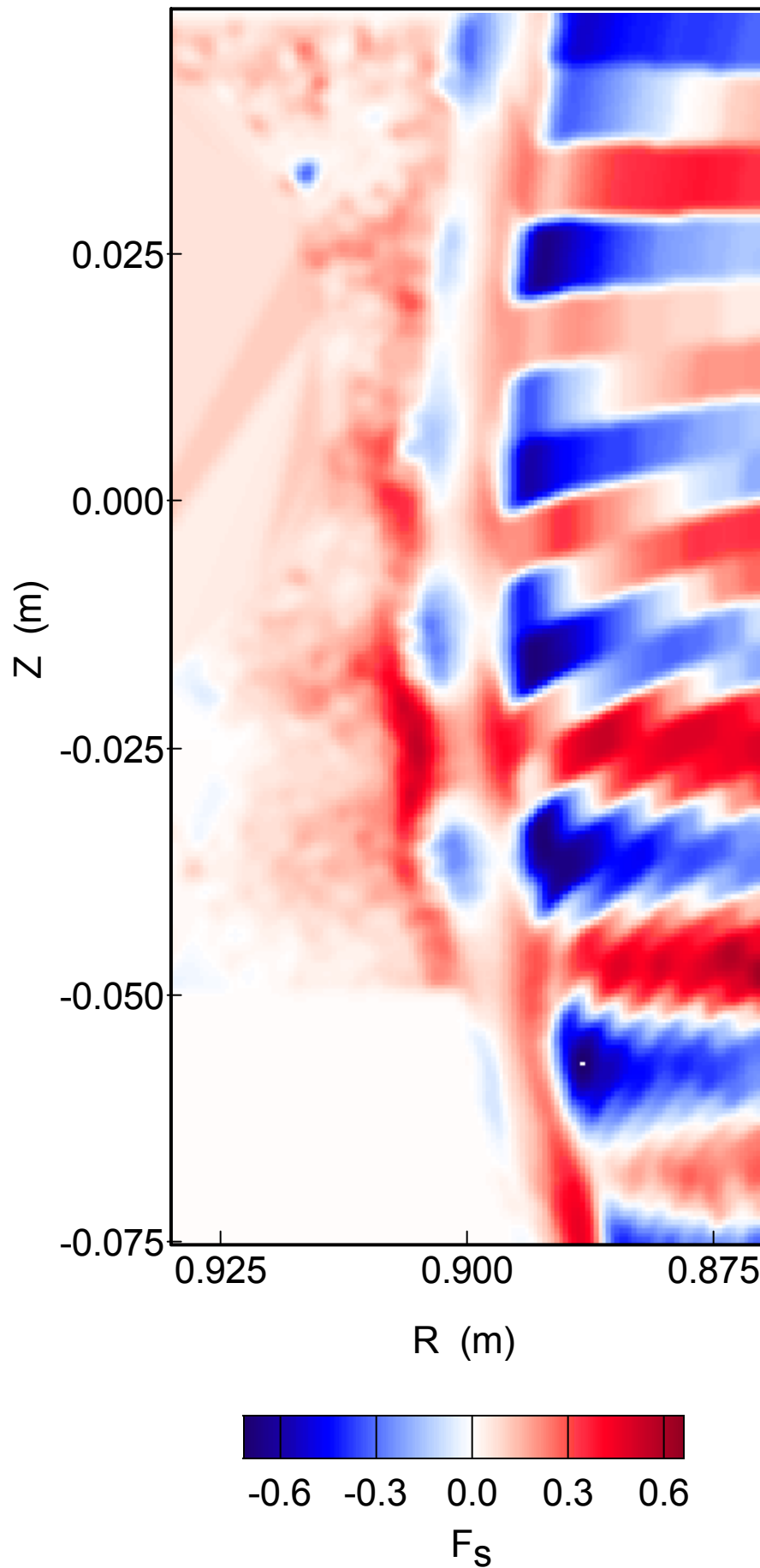
- Main observations:

1.  $|F_s| \gtrsim 0.5$  in many places  
 $\Rightarrow$  too large to ignore in GPI analysis.
2. Most of  $F_s$  due to molecules,  
  - Analogous quantity based on atoms only  $\leq 0.2$ .

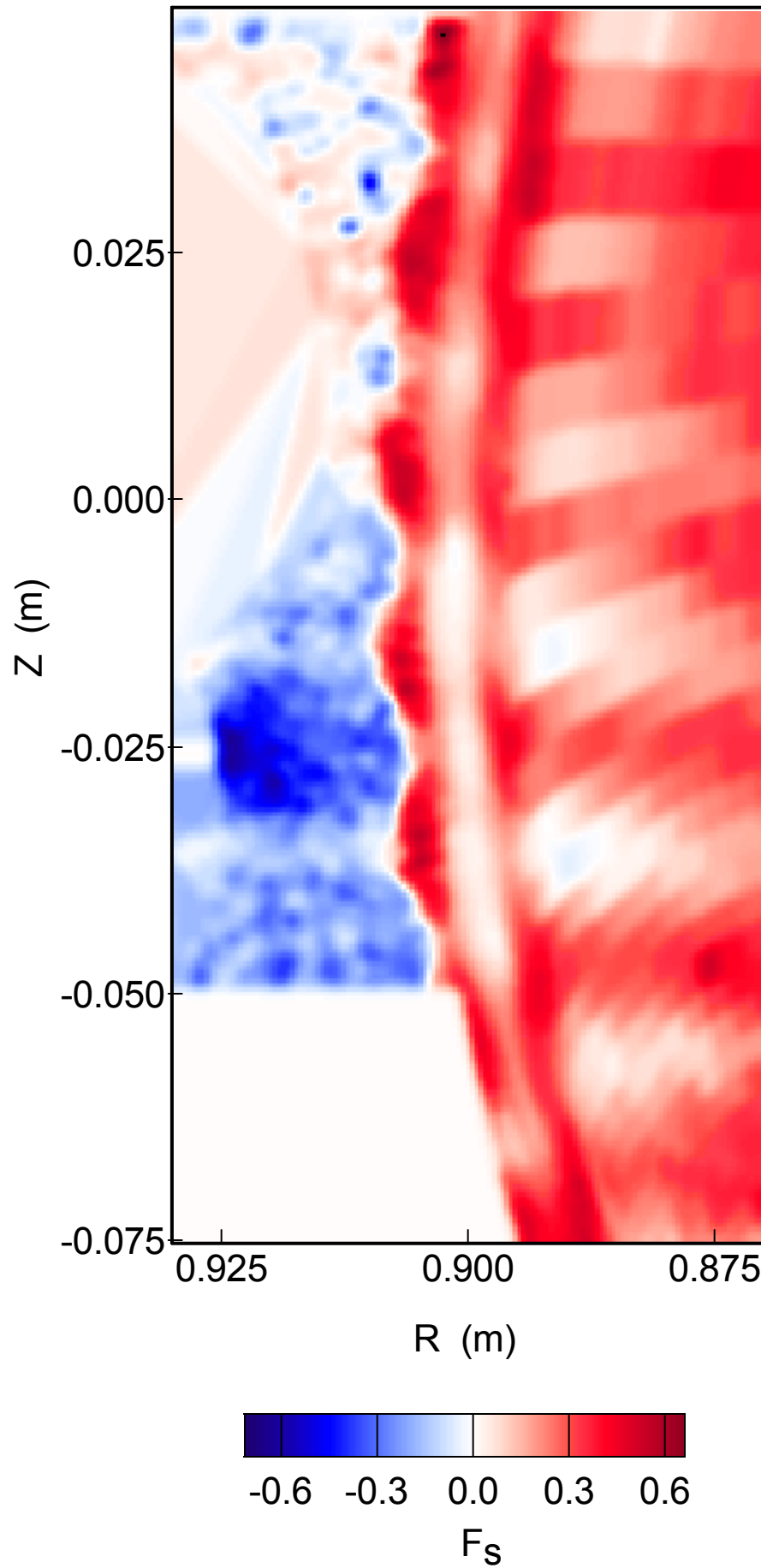
- To understand  $F_s$  look at radial slices,

- $Z = -0.034$ : peak in  $n'_e/n_e$ ,
- $Z = -0.025$ : at nozzle & a minimum in  $n'_e/n_e$ .
- Compare with  $1 - n'_e/n_e$ ,
  - \*  $1 - n'_e/n_e < 0 \Rightarrow$  local  $n_e >$  unperturbed value,
  - \*  $1 - n'_e/n_e > 0 \Rightarrow$  local  $n_e <$  unperturbed value,
  - \*  $T_e$  perturbation differs at edges.
- $F_s < 0 \Rightarrow n_j$  locally reduced,
  - \*  $F_s$  drops are in “shadows” of largest  $n'_e/n_e$ .
- $F_s > 0 \Rightarrow n_j$  locally increased,
  - \*  $F_s > 0$  at  $Z = -0.025$  since  $n_e$  modulation near min.,
  - \* Not so in perturbed  $T_e$  case due to smaller dissociation rate & strong  $T_e$  dependence of  $f_j$ .

# Shadow Fraction with Density Perturbation

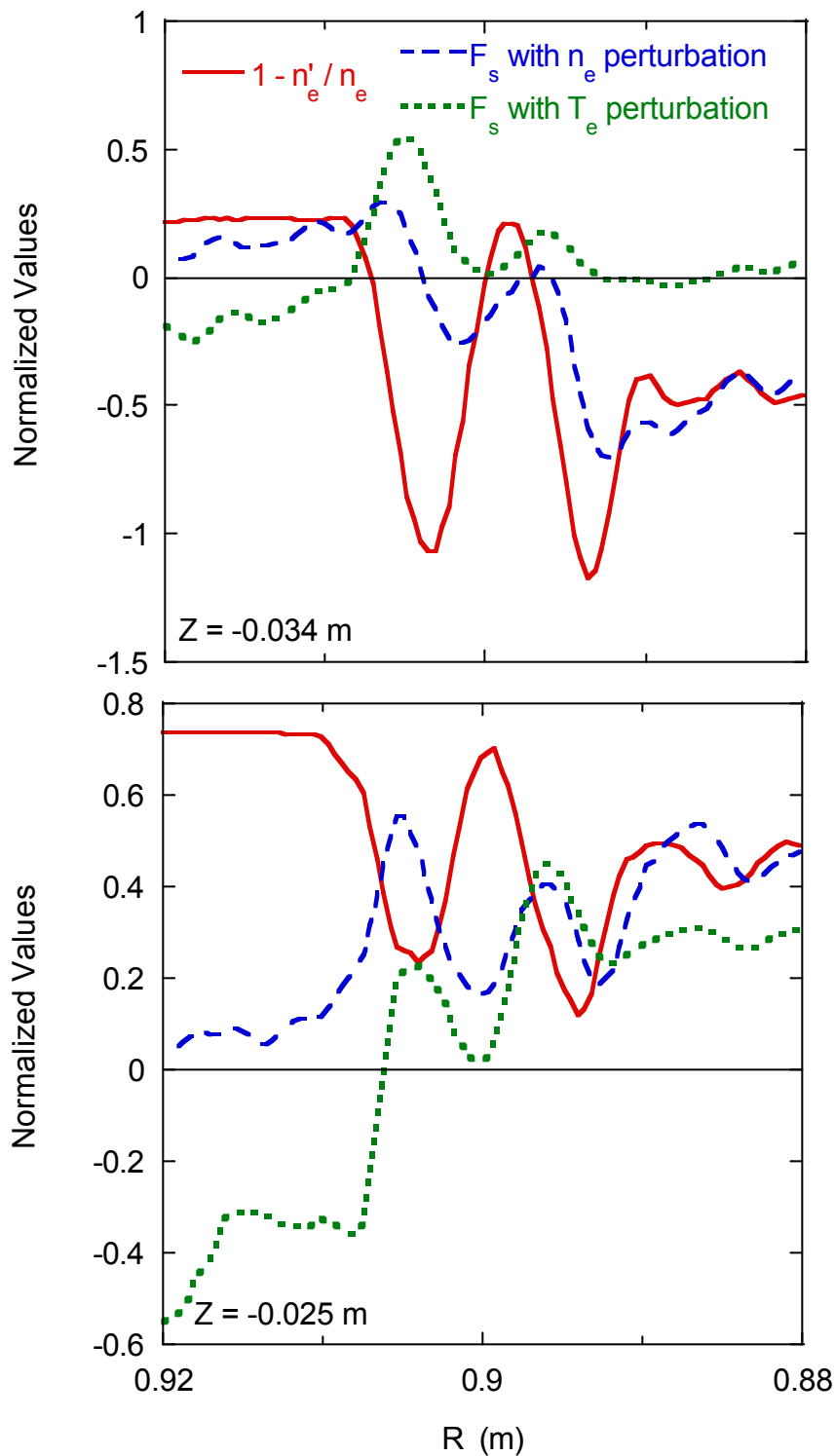


# Shadow Fraction with Temperature Perturbation



# Shadow Fraction Significant

## Radial Slices



# CONCLUSIONS

- DEGAS 2 simulations show that spatial variation of  $D_\alpha$  emission reflects that of  $n_e, T_e$  turbulence.
- But,  $n_e, T_e$  dependence of emission rate complicated,
  - $\Rightarrow$  no simple scheme to get plasma fluctuations.
- Contributions from molecules significant,
  - Further complicating  $n_e, T_e$  dependence,
  - Densities significantly affected by perturbation.
- $\Rightarrow$  will need neutral transport code to interpret GPI,
  - Must do careful benchmarks first,
  - To verify these conclusions,
  - Validate atomic & molecular physics models.